Background



Figure 3.1 Model of an automobile suspension with input and output signals.

Here is an example of a second order system from EAS 206.





We can write the equation in the form

$$m\frac{d^2x}{dt^2} + c\frac{dx}{dt} + kx = F(t)$$

consider first the simpler case where c = 0 (no damping) and F(t) = 0. The equation becomes

$$m\frac{d^2x}{dt^2} + kx = 0$$

which has the solution

$$y = C\sin(\omega_n t)$$

The mass will oscillate sinusoidally and the oscillation will continue forever at the undamped natural frequency ω_n

Recognizing the periodic nature of the solution, it is convenient to rewrite the equation in the form

$$\frac{1}{\omega_n^2} \frac{d^2 y}{dt^2} + \frac{2\zeta}{\omega_n} \frac{dy}{dt} + y = KF(t)$$
(3.13)

where ω_n is the natural frequency and ζ (zeta) is the damping ratio.

 $\omega_n = \sqrt{\frac{k}{m}}$ = natural frequency of the system $\zeta = \frac{c}{2\sqrt{km}}$ = damping ratio of the system

Homogeneous Solution

The form of the homogeneous solution depends on the roots of the characteristic equation

$$\frac{1}{\omega_n^2}\lambda^2 + \frac{2\zeta}{\omega_n}\lambda + 1 = 0$$

The quadratic equation has two roots,

$$\lambda_{1,2} = -\zeta \omega_n \pm \omega_n \sqrt{\zeta^2 - 1}$$

Depending on the value of ζ , three forms of the homogeneous solution are possible:

 $0 < \zeta < 1$ (under damped system solution)

$$y_{h}(t) = Ce^{-\zeta \omega_{n} t} \sin \omega_{n} \sqrt{1 - \zeta^{2}} t + \Theta$$
(3.14a)

 $\zeta = 1$ (critically damped system solution)

$$y_h(t) = C_1 e^{\lambda_1 t} + C_2 t e^{\lambda_2 t}$$
 (3.14b)

 $\zeta > 1$ (over damped system solution)

$$y_h(t) = C_1 e^{\lambda_1 t} + C_2 e^{\lambda_2 t}$$
 (3.14c)

The particular solution will depend on the forcing function F(t).

Step Function Input

For an underdamped system, $0 < \zeta < 1$, F(t) = AU(t), the solution of equation (3.13) is:

$$y(t) = KA - KAe^{-\xi\omega_{n}t} \left[\frac{\zeta}{1 - \zeta^{2}} \sin \omega_{n} \sqrt{1 - \zeta^{2}}t + \cos \omega_{n} \sqrt{1 - \zeta^{2}}t \right]_{(3.15a)}$$

For a critically damped system the solution is:

$$y(t) = KA - KA(1 + \omega_n t)e^{-\omega_n t}$$
(3.15b)

For an overdamped system the solution is:

$$y(t) = KA - KA \left[\frac{\zeta + \sqrt{\zeta^2 - 1}}{2\sqrt{\zeta^2 - 1}} e^{-\zeta + \sqrt{\zeta^2 - 1} \omega_n t} + \frac{\zeta - \sqrt{\zeta^2 - 1}}{2\sqrt{\zeta^2 - 1}} e^{-\zeta - \sqrt{\zeta^2 - 1} \omega_n t} \right]$$
(3.15c)

FIGURE 3.13 Second-order system response to a step function input.



For underdamped systems, the output oscillates at the ringing frequency ω_d

$$T_d = \frac{2\pi}{\omega_d} = \frac{1}{f_d} \tag{3.16}$$

$$\omega_d = \omega_n \sqrt{1 - \zeta^2} \tag{3.17}$$

Remember $f = \omega/2\pi$

Rise Time

By definition it is the time required for the system to achieve a value of 90% of the step input. The rise time is decreased by decreasing the damping (see figure 3.13 above). Obviously there is a tradeoff between fast response and ringing in a second order system.

Settling Time

The settling time is defined as the time required for the system to settle to within $\pm 10\%$ of the steady state value.

A damping ratio, ζ , of 0.7 offers a good compromise between rise time and settling time. Most dynamic response measurement systems are designed such that the damping ratio is between 0.6 and 0.8

Frequency Response

If $F(t) = A \sin \omega t$, the solution is given by

$$y(t) = y_{h} + \frac{KA \sin[\omega t + \phi(\omega)]}{\{[1 - (\omega/\omega_{n})^{2}]^{2} + (2\zeta\omega/\omega_{n})^{2}\}^{1/2}}$$
(3.18)

The first term is a transient which will eventually die out - the steady-state response can be written in the form

$$y(t \to \infty) = M(\omega) KA \sin[\omega t + \phi(\omega)]$$
 (3.20)

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where

$$\phi(\omega) = \tan^{-1} \left(-\frac{2\zeta \, \omega/\omega_n}{1 - (\omega/\omega_n)^2} \right)$$
(3.19)

$$M(\omega) = \frac{B}{KA} = \frac{1}{\{[1 - (\omega/\omega_n)^2]^2 + (2\zeta\omega/\omega_n)^2\}^{1/2}}$$
(3.21)

EXAMPLE 3.8

Find the rise and settling time and damped natural frequency of the second order system step input response in figure 3.15.

From the figure y_0 is determined to be 1 V (note error in y axis label) and $y_{\infty} = 2$ V.

Therefore 90% of $(y_{\infty}-y_0)$ is 0.9 V. The rise time can be determined by locating the point at which the system response reaches $y_0 + 0.9V = 1.9$ V. The settling time is determined by locating the point at which the system remains with the range of $y_{\infty} \pm 10\%(y_{\infty}-y_0) = 2 \pm 0.1V$

The damped natural frequency or ringing frequency is found by determining the period of the oscillation, T_d , and recalling the relation between period in seconds, frequency in cycles per second and the conversion to circular frequency, radians/second. From the graph T_d is found to be 13 ms. Therefore f_d = 1/13 ms = $\omega_d/2\pi$.



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Resonance Frequency

Underdamped second order systems may resonate or oscillate at a greater magnitude than the input, $M(\omega) > 1$.

Resonance Band: is the frequency range over which $M(\omega) > 1$.

Resonance Frequency: is defined as
$$\omega_R = \omega_n \sqrt{1 - 2\zeta^2}$$

Systems with a damping ratio greater than ζ > 0.707 do not resonate.





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EXAMPLE 3.10

Select and appropriate accelerometer natural frequency and damping ratio to measure frequencies below 100 Hz (628 rad/s) and maintain a dynamic error of ±5%. (i.e. $M(\omega) \ge 0.95$)

Assume that the accelerometer has a damping ratio of most dynamic sensors of 0.7. Therefore using equation 3.21 one could solve for ω_n or plot the equation with $\zeta = 0.7$ and find the frequency range over which $1.05 \le M(\omega) \ge 0.95$.



Figure 3.18 Magnitude ratio for second-order system with $\zeta = 0.7$ for Example 3.10.



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