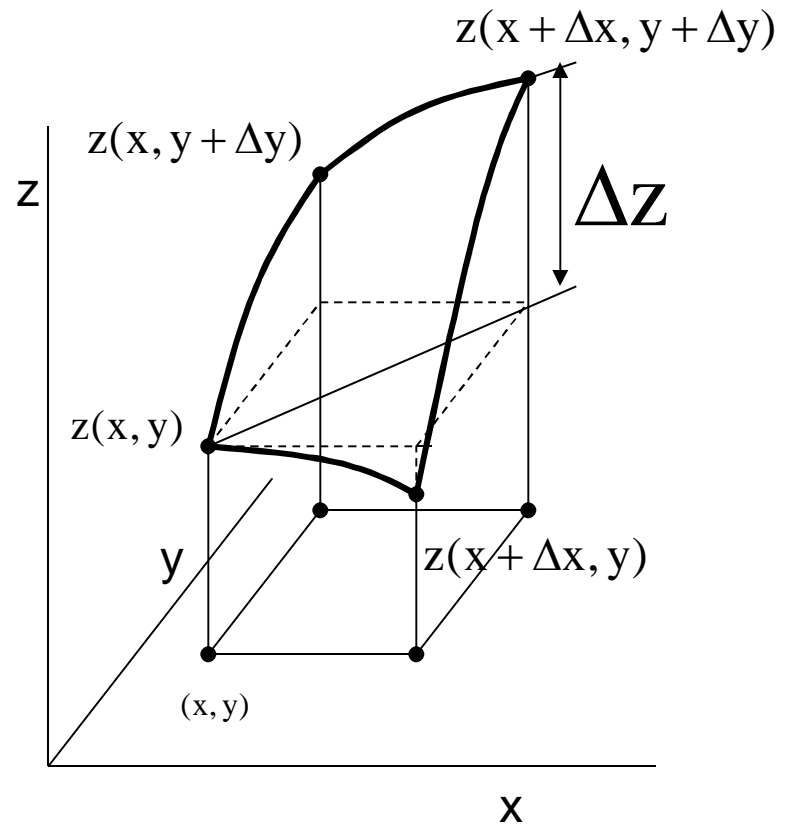
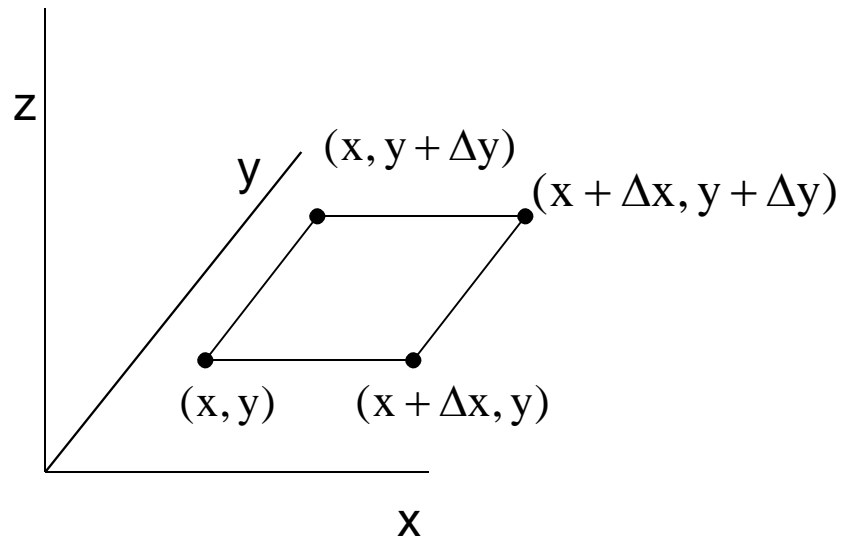


Exact Differential



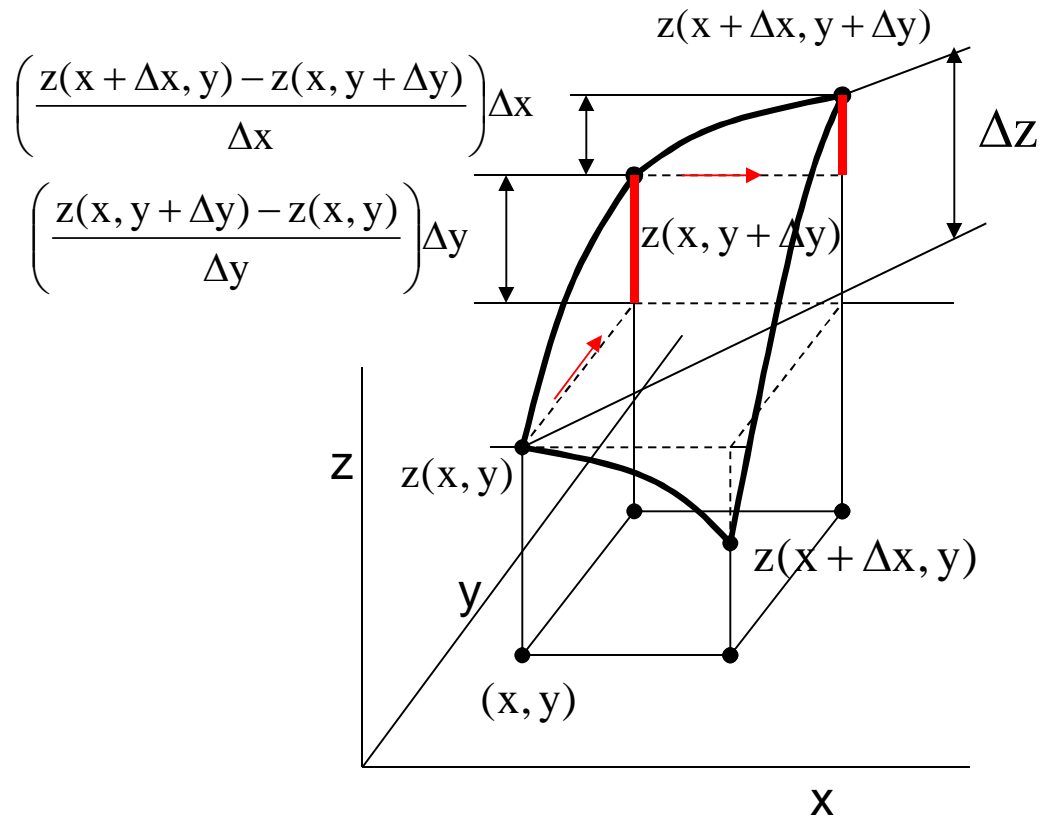
Exact Differential

Theorem :

If the function $z = f(x, y)$ has continuous first partial derivatives in the domain, then the function z has a differential

$$dz = \frac{\partial z}{\partial x} dx + \frac{\partial z}{\partial y} dy$$

at every point in the domain.



$$dz = \left(\frac{z(x + \Delta x, y) - z(x, y + \Delta y)}{\Delta x} \right) \Delta x + \left(\frac{z(x, y + \Delta y) - z(x, y)}{\Delta y} \right) \Delta y$$

$$\lim_{\Delta x \rightarrow 0} \frac{z(x + \Delta x, y) - z(x, y)}{\Delta x} = \left(\frac{\partial z}{\partial x} \right)_{y=\text{const}}$$

$$dz = \left(\frac{\partial z}{\partial x} \right)_{y=\text{const}} dx + \left(\frac{\partial z}{\partial y} \right)_{x=\text{const}} dy \quad (12-3)$$

$$dz = \left(\frac{\partial z}{\partial x} \right)_{y=\text{const}} dx + \left(\frac{\partial z}{\partial y} \right)_{x=\text{const}} dy$$

differentiating the partial derivatives again,

$$\frac{\partial}{\partial y} \left(\frac{\partial z}{\partial x} \right) = \frac{\partial}{\partial x} \left(\frac{\partial z}{\partial y} \right) \text{ since order is immaterial}$$

Since $x = x(y, z)$

$$dx = \left(\frac{\partial x}{\partial y} \right) dy + \left(\frac{\partial x}{\partial z} \right) dz \quad \text{and} \quad dy = \left(\frac{\partial y}{\partial x} \right) dx + \left(\frac{\partial y}{\partial z} \right) dz$$

substituting into dz,

$$\left(\frac{\partial x}{\partial y} \right)_z \left(\frac{\partial y}{\partial z} \right)_x = 1$$

$$\text{and} \left(\frac{\partial x}{\partial y} \right)_z \left(\frac{\partial y}{\partial z} \right)_x \left(\frac{\partial z}{\partial x} \right)_y = -1 \quad (12-19)$$

Maxwell Relation Derivation

$$q = du + pdv \quad \text{First Law}$$

$$q = Tds \quad \text{Second Law}$$

Combining,

$$du = [T]ds + [-p]dv$$

Since u is a property, independent of path and a function of any two other properties, $u = u(s, v)$, du is an exact differential.

$$du = \left[\left(\frac{\partial u}{\partial s} \right)_v \right] ds + \left[\left(\frac{\partial u}{\partial v} \right)_s \right] dv$$

comparing the quantities in bracket,

$$T = \left(\frac{\partial u}{\partial s} \right)_v \quad \text{and} \quad -p = \left(\frac{\partial u}{\partial v} \right)_s$$

differentiating again,

$$\left(\frac{\partial T}{\partial v} \right)_s = \frac{\partial}{\partial v} \left(\frac{\partial u}{\partial s} \right)_v \quad \text{and}$$

$$\left(-\frac{\partial p}{\partial s} \right)_v = \frac{\partial}{\partial s} \left(\frac{\partial u}{\partial v} \right)_s$$

order does not affect the result of differentiation and,

$$\frac{\partial}{\partial v} \left(\frac{\partial u}{\partial s} \right)_v = \frac{\partial}{\partial s} \left(\frac{\partial u}{\partial v} \right)_s$$

substituting,

$$\left(\frac{\partial T}{\partial v} \right)_s = - \left(\frac{\partial p}{\partial s} \right)_v \quad (12-16)$$

Using $du = Tds - pdv$ (12-10)

three relationships were derived

$$T = \left(\frac{\partial u}{\partial s} \right)_v$$

$$-p = \left(\frac{\partial u}{\partial v} \right)_s$$

$$\left(\frac{\partial T}{\partial v} \right)_s = - \left(\frac{\partial p}{\partial s} \right)_v \quad (12-16)$$

Using three additional equations,

First Law + h definition

$$dh = Tds + vdp \quad (12-11)$$

Helmholtz function

a property composed of properties

$$a = u - Ts \quad (12-12)$$

GibbvsFunction (Availability)

a property composed of properties

$$g = h - Ts \quad (12-13)$$

using this same technique 9 additional

relationships can be derived.

Maxwell Relations

from $u = u(s, v)$

$$T = \left(\frac{\partial u}{\partial s} \right)_v \quad -p = \left(\frac{\partial u}{\partial v} \right)_s$$

from $h = h(s, p)$

$$T = \left(\frac{\partial h}{\partial s} \right)_p \quad v = \left(\frac{\partial h}{\partial p} \right)_s$$

$$\left(\frac{\partial T}{\partial v} \right)_s = - \left(\frac{\partial p}{\partial s} \right)_v \quad (12-16)$$

$$\left(\frac{\partial T}{\partial p} \right)_s = - \left(\frac{\partial v}{\partial s} \right)_p \quad (12-17)$$

$$\left(\frac{\partial p}{\partial T} \right)_v = \left(\frac{\partial s}{\partial v} \right)_T \quad (12-18)$$

$$\left(\frac{\partial v}{\partial T} \right)_p = - \left(\frac{\partial s}{\partial p} \right)_T \quad (12-19)$$

Additional Relations

Helmholtz Function $a = u - Ts$

from $a = a(v, T)$

$$-p = \left(\frac{\partial a}{\partial v} \right)_T \quad v = \left(\frac{\partial a}{\partial T} \right)_p$$

Gibbs Function $g = h - Ts$

from $g = g(T, p)$

$$v = \left(\frac{\partial g}{\partial p} \right)_T \quad v = \left(\frac{\partial g}{\partial T} \right)_p$$

$$\left(\frac{\partial u}{\partial s} \right)_v = \left(\frac{\partial h}{\partial s} \right)_p$$

$$\left(\frac{\partial h}{\partial p} \right)_s = \left(\frac{\partial g}{\partial p} \right)_T$$

$$\left(\frac{\partial u}{\partial v} \right)_s = - \left(\frac{\partial a}{\partial v} \right)_T$$

$$\left(\frac{\partial a}{\partial T} \right)_v = - \left(\frac{\partial g}{\partial T} \right)_p$$

$$\left(\frac{\partial v}{\partial T} \right)_p = - \left(\frac{\partial s}{\partial p} \right)_T$$

Given that enthalpy of a gas is known everywhere as a function of pressure and entropy, $h = h(s, p)$ find t, v, u, g and a as functions of s and p .

$$dh = \left(\frac{\partial h}{\partial s} \right)_p ds + \left(\frac{\partial h}{\partial p} \right)_s dp \quad \text{exact differential of the property } h$$

$$dh = Tds + vdp \quad \text{First Law and definition of enthalpy}$$

$$\text{a) comparing equations, } T(s, p) = \left(\frac{\partial h}{\partial s} \right)_p$$

$$\text{b) comparing equations, } v(s, p) = \left(\frac{\partial h}{\partial p} \right)_s$$

$$\text{c) } u = h - pv \quad \text{definition of } h$$

$$u(s, p) = h(s, p) - p \left(\frac{\partial h}{\partial p} \right)_s$$

$$\text{d) } a(s, p) = u - Ts$$

$$a(s, p) = h(s, p) - p \left(\frac{\partial h}{\partial p} \right)_s - \left(\frac{\partial h}{\partial s} \right)_p s$$

$$\text{e) } g(s, p) = h - Ts$$

$$g(s, p) = h(s, p) - \left(\frac{\partial h}{\partial s} \right)_p s$$

$$z = z(x, y)$$

$$\left(\frac{\partial x}{\partial y}\right)_z \left(\frac{\partial y}{\partial z}\right)_x \left(\frac{\partial z}{\partial x}\right)_y = -1$$

$$h = h(p, T)$$

$$\left(\frac{\partial p}{\partial T}\right)_h \left(\frac{\partial T}{\partial h}\right)_p \left(\frac{\partial h}{\partial p}\right)_T = -1$$

$$c_p = \left(\frac{\partial h}{\partial T}\right)_p \text{ Specific Heat}$$

at constant pressure

$$JT = \left(\frac{\partial T}{\partial p}\right)_h \text{ Joule Thompson}$$

coefficient

$$c = \left(\frac{\partial h}{\partial p}\right)_T \text{ constant temperature}$$

coefficient

$$\frac{1}{JT} \times \frac{1}{c_p} \times c = -1$$

$$c = -JT \times c_p$$

$$\left(\frac{\partial T}{\partial p} \right)_h = \text{JT Coefficient}$$

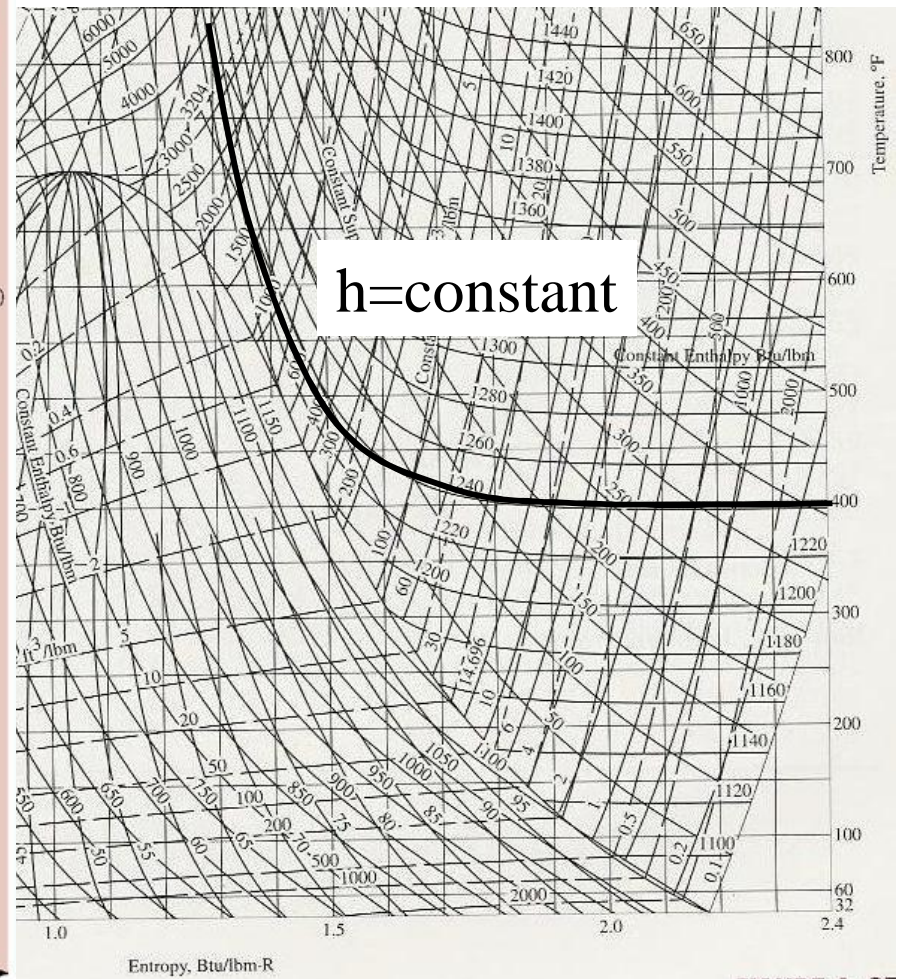
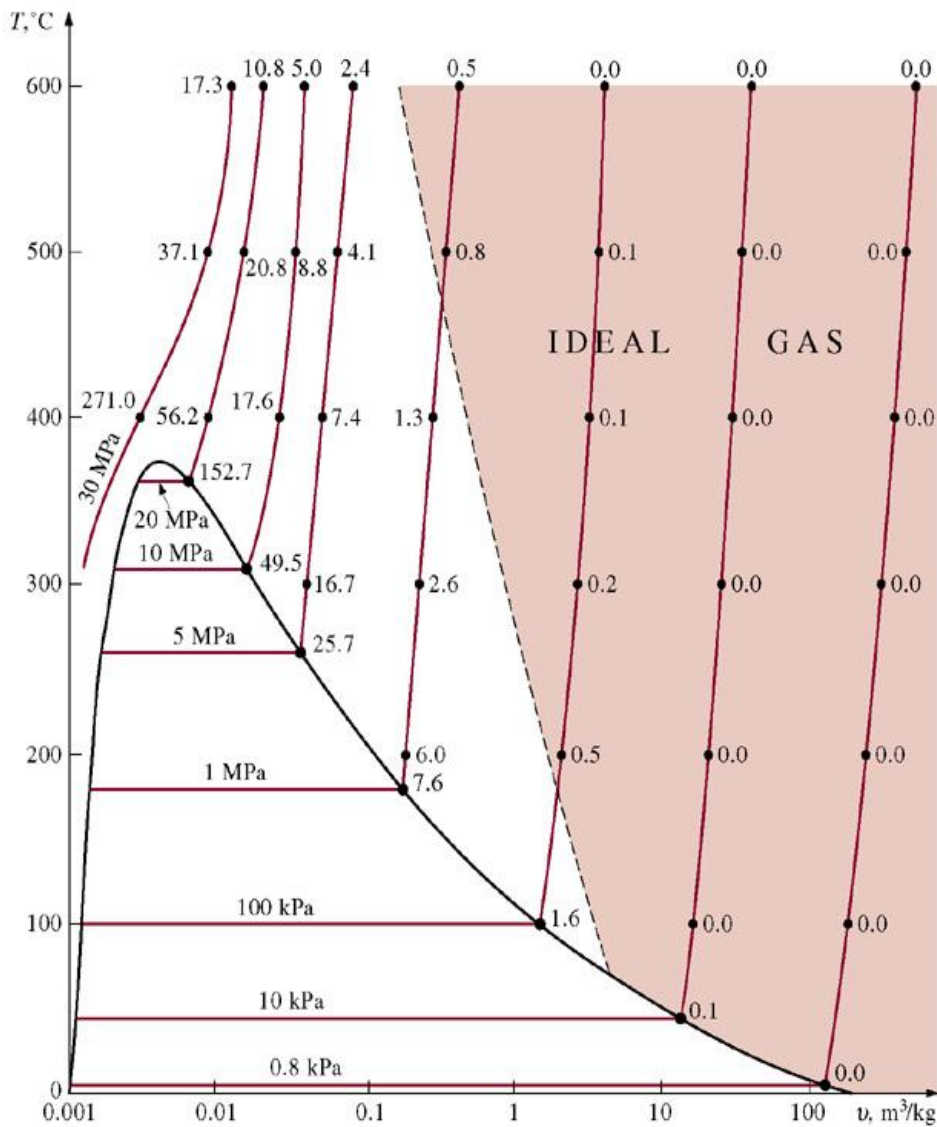


FIGURE A-9E

T-s diagram for water. [Source: Joseph H. Keenan, Frederick G. Keyes, Philip G. Hill, and Joan G. Moore, *Steam Tables* (New York: John Wiley & Sons, 1969).]

Clapeyron Relation

Using the Maxwell Relation,

$$\left(\frac{\partial s}{\partial v}\right)_T = \left(\frac{\partial p}{\partial T}\right)_v \quad (12.18) \quad \left(\frac{\partial h}{\partial s}\right)_p = T$$

for

a two phase mixture

of a pure substance

$$\left(\frac{\partial s}{\partial v}\right)_T = \left(\frac{s_{fg}}{v_{fg}}\right)_T = \left(\frac{\partial p}{\partial T}\right)_v$$

2 phase

\Rightarrow p is a function only of T

$$\left(\frac{s_{fg}}{v_{fg}}\right)_T = \frac{dp}{dT}$$

$$\frac{h_{fg}}{s_{fg}} = T$$

$$s_{fg} = \frac{h_{fg}}{T}$$

substituting gives the Clapeyron Relation

$$\left(\frac{dp}{dT}\right)_{\text{sat}} = \frac{h_{fg}}{T v_{fg}} \quad (12.22)$$

relates h_{fg} and v_{fg} simply to p, v, T data

PRINCIPAL OF CORRESPONDING STATES

COMPRESSIBILITY FACTOR Z

Z is **about** the same for all gases at the same reduced temperature and the same reduced pressure where:

$$Z = \frac{pv}{RT}$$

$$Z = \frac{pV}{mRT}$$

$$P_R = \frac{p}{p_{\text{critical}}} \quad (3-22)$$

$$T_R = \frac{T}{T_{\text{critical}}}$$

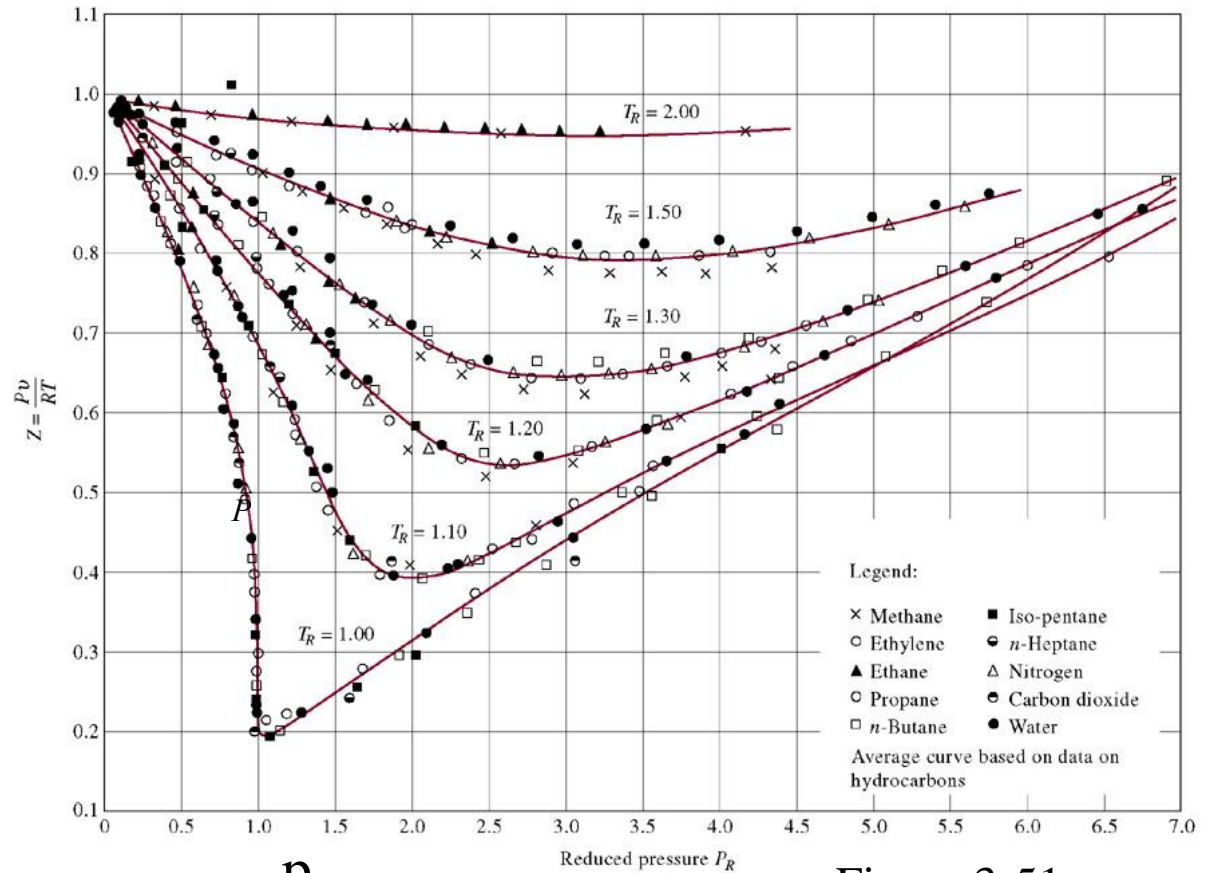


Figure 3-51

VAN DER WAALS EQUATION OF STATE - 1873

$$\left(p + \frac{a}{v^2} \right) (v - b) = RT \quad (3-22)$$

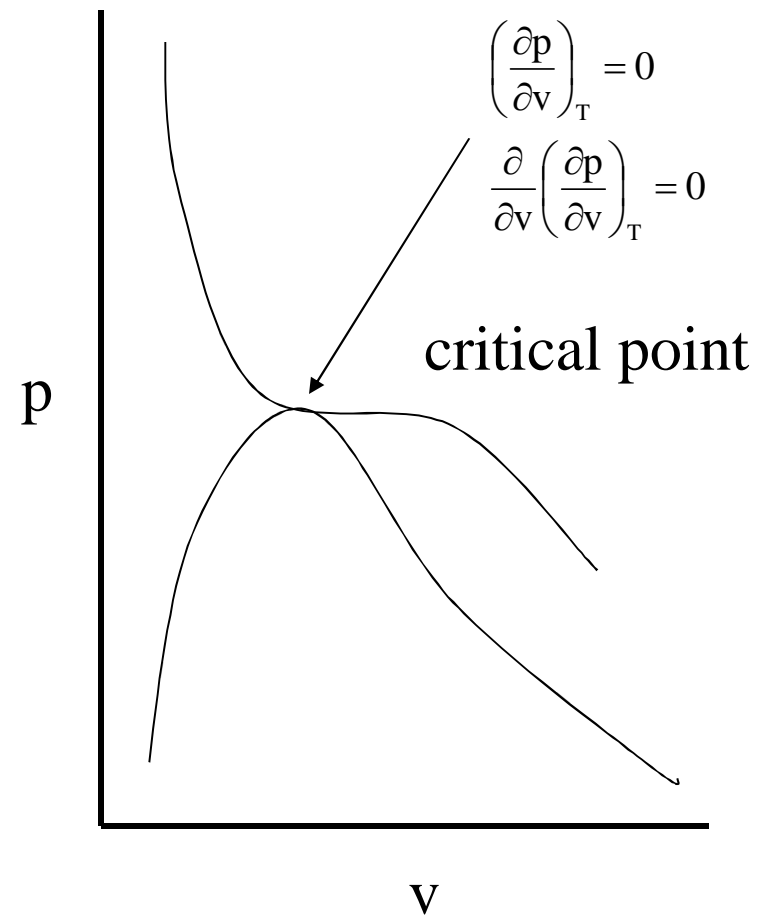
$\frac{a}{v^2}$ – intermolecular forces

b – volume of gas molecules

$$p = \frac{RT_{\text{critical}}}{v_{\text{critical}} - b} - \frac{a}{v_{\text{critical}}^2}$$

$$\left(\frac{\partial p}{\partial v} \right)_{T_{\text{critical}}} = 0 \quad \left(\frac{\partial^2 p}{\partial v^2} \right)_{T_{\text{critical}}} = 0$$

$$a = \frac{27 R^2 T_{\text{critical}}^2}{64 p_{\text{critical}}} \quad b = \frac{R T_{\text{critical}}}{8 p_{\text{critical}}}$$



Equations of State

Ideal Gas Law	$pV = RT$	Beatti Bridgman 1928 (3-24) 5 constants Table 3-4 coefficients
Compressibility	$pV = ZRT$ Figure 3-51	Redlich Kwong 1949
Variable Specific Heat	$c_p(T)$ $h = \int c_p(T) dT$ Table A-17 to 23 Table A-17E to 23E	Benedict Webb Rubin (3-26) 8 constants Table A-24 coefficients
Van der Waals Equation	1873 (3-22) $\left(p + \frac{a}{v^2}\right)(v - b) = RT$	Virial Equation $p = \frac{RT}{v} + \frac{a(T)}{v^2} + \frac{b(T)}{v^3} + \dots \quad (3-27)$ Spread Sheet World air equation - 32 constants Many, many others.

Thermodynamic Properties are Measurements

p, T, v, u, h, s - measure directly

-measure by change $\left(\frac{\partial s}{\partial v}\right)_T = \left(\frac{\partial p}{\partial T}\right)_v$

