

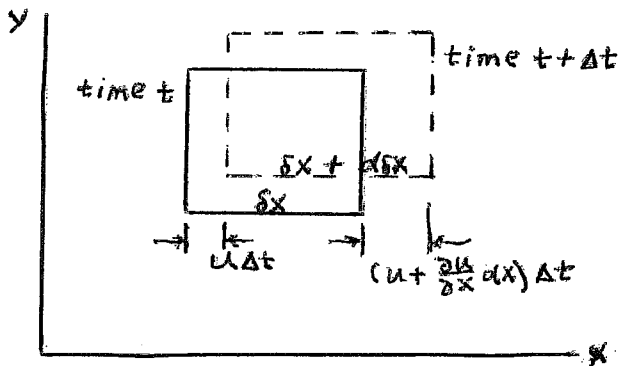
Conservation of matter

$$\frac{D\delta mass}{Dt} = 0 \quad \text{where } \delta mass = \rho \delta \nabla$$

$$\frac{D\rho \delta \nabla}{Dt} = \rho \frac{D\delta \nabla}{Dt} + \delta \nabla \frac{D\rho}{Dt} = 0$$

$$\frac{D\rho}{Dt} + \rho \frac{1}{\delta \nabla} \frac{D\delta \nabla}{Dt} = 0$$

$(D\delta \nabla / Dt) / \delta \nabla = \text{volume strain rate}$



New horizontal length of the rectangle

$$\delta x + d\delta x = \delta x + \left(u + \frac{\partial u}{\partial x} dx\right) dt - u dt \quad \text{and} \quad d\delta x = \frac{\partial u}{\partial x} dx dt$$

$$\text{Strain rate for the } \delta x \text{ element} = \frac{1}{\delta x} \frac{d\delta x}{dt} = \frac{\partial u}{\partial x}$$

$$\text{For } \delta y \text{ and } \delta z: \frac{1}{\delta y} \frac{d\delta y}{dt} = \frac{\partial v}{\partial y} \quad \text{and} \quad \frac{1}{\delta z} \frac{d\delta z}{dt} = \frac{\partial w}{\partial z}$$

Displaced volume

$$\begin{aligned}\delta\forall + d\delta\forall &= (\delta x + d\delta x)(\delta y + d\delta y)(\delta z + d\delta z) \\ &\cong \delta x \delta y \delta z + \delta y \delta z d\delta x + \delta x \delta z d\delta y + \delta x \delta y d\delta z\end{aligned}$$

$$d\delta\forall \cong \delta y \delta z d\delta x + \delta x \delta z d\delta y + \delta x \delta y d\delta z$$

$$\frac{1}{\delta\forall} \frac{d\delta\forall}{dt} = \frac{1}{\delta x} \frac{d\delta x}{dt} + \frac{1}{\delta y} \frac{d\delta y}{dt} + \frac{1}{\delta z} \frac{d\delta z}{dt} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}$$

Continuity equation

$$\frac{D\rho}{Dt} + \rho \frac{1}{\delta\forall} \frac{D\delta\forall}{Dt} = \frac{D\rho}{Dt} + \rho \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) = 0$$

= Change in density + change in volume

Lagrangian derivative for density

$$\frac{D\rho}{Dt} = \frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} + w \frac{\partial \rho}{\partial z}$$

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} + w \frac{\partial \rho}{\partial z} + \rho \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) = 0$$

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho u}{\partial x} + \frac{\partial \rho v}{\partial y} + \frac{\partial \rho w}{\partial z} = 0$$

Same as obtained via the control volume analysis.

For incompressible flow

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 = \text{volume stain rate}$$

Example For a steady velocity field, $u=ax$ and $v=by$, determine the density variation in terms of the Lagrangian and Eulerian formulations.

$$\text{Volume strain rate} = \partial u / \partial x + \partial v / \partial y = a + b$$

$$\frac{D\rho}{Dt} + \rho \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) = 0$$

$$\frac{d\rho}{dt} + (a+b)\rho = 0 \text{ which has the solution } \rho = \rho_0 e^{-(a+b)t}$$

Special case: $b=-a$, $\rho = \rho_0$, and the fluid is incompressible.

$$\text{Particle trajectories: } \frac{dX}{dt} = u = aX \text{ and } \frac{dY}{dt} = v = bY$$

$$\text{Solutions: } X = X_0 e^{at} \text{ and } Y = Y_0 e^{bt}$$

$$Y = Y_0 (X / X_0)^{b/a}$$

Particle path in terms of field coordinates

$$y = y_0 (x / x_0)^{b/a}$$

With the equations for X,Y; $\rho = \rho_0 e^{-(a+b)t}$ becomes

$$\rho = \rho_0 \frac{X_0 Y_0}{XY} \text{ and } \rho = \rho_0 \frac{x_0 y_0}{xy}$$

Eulerian formulation

$$u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} + \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \rho = 0$$

$$ax \left(-\rho_0 \frac{x_0 y_0}{x^2 y} \right) + by \left(-\rho_0 \frac{x_0 y_0}{xy^2} \right) + (a+b) \rho_0 \frac{x_0 y_0}{xy} \equiv 0$$

which is satisfied.

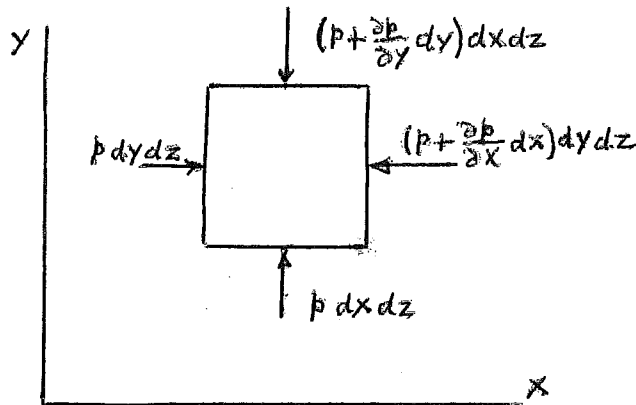
Momentum Equation

Newton's Second Law for a system of mass $dm = \rho dV$

$$\frac{Ddm\vec{V}}{Dt} = dm \frac{D\vec{V}}{Dt} = \rho dV \frac{D\vec{V}}{Dt} = \sum d\vec{F} \quad 3.3.1$$

$d\vec{F}$ = pressure, gravity or friction force acting on the system

Pressure Forces on System of rectangular shape



Net pressure force

$$\begin{aligned} d\vec{F}_p = & [p dy dz - (p + \frac{\partial p}{\partial x} dx) dy dz] \hat{i} \\ & + [p dx dz - (p + \frac{\partial p}{\partial y} dy) dx dz] \hat{j} \\ & + [p dx dy - (p + \frac{\partial p}{\partial z} dz) dx dy] \hat{k} \end{aligned}$$

Then
$$d\vec{F}_p = -(\frac{\partial p}{\partial x} \hat{i} + \frac{\partial p}{\partial y} \hat{j} + \frac{\partial p}{\partial z} \hat{k}) dx dy dz = -\nabla p dV$$

$$\nabla p = \hat{i} \partial p / \partial x + \hat{j} \partial p / \partial y + \hat{k} \partial p / \partial z = \text{gradient operator}$$

$$-\nabla p = \text{net pressure force per unit volume}$$

Gravity force = body force acts on the mass of the system

Pressure force = surface force

Magnitude of the gravity force is dmg

Direction toward the center of the earth given by \vec{g}

$$d\vec{F}_g = dm\vec{g} = \rho\vec{g}d\forall$$

Friction force = surface stresses, both normal and tangential

$$d\vec{F}_f = \vec{f}d\forall$$

\vec{f} = the net friction force per unit volume

Equation of motion for unsteady, compressible, frictional flow

$$\rho \frac{D\vec{V}}{Dt} = -\nabla p + \rho\vec{g} + \vec{f}$$

Each term then has dimensions of force per unit volume.

$$\rho \left(\frac{\partial \vec{V}}{\partial t} + u \frac{\partial \vec{V}}{\partial x} + v \frac{\partial \vec{V}}{\partial y} + w \frac{\partial \vec{V}}{\partial z} \right) = -\nabla p + \rho\vec{g} + \vec{f}$$

$$\vec{V} = \vec{V}(x, y, z, t)$$

Along with the continuity of equation there are four equations to determine u, v, w and p as functions of x, y, z, t .